# Molecular Structure of the GeH<sub>2</sub>…OH<sub>2</sub> Complex

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Hartree–Fock and post-Hartree–Fock (density functional DFT, many-body perturbation theory MP2 and MP4(SDTQ), and coupled-cluster CCSD, and CCSD[T]) quantum-mechanical methods were applied to establish molecular structures and interaction energies of the GeH<sub>2</sub>···OH<sub>2</sub> complex. Molecular geometries of nonplanar and planar conformers were optimized at the Hartree–Fock, DFT, and MP2 levels of theory using triple- $\zeta$ -quality (TZP) augmented by sets of polarization functions (TZP(2d,2p) and TZP(2df,2pd)) basis sets. The optimized nonplanar structures correspond to minimum-energy species at all applied levels of theory. Calculated interaction energies (corrected for the basis set superposition error, CP) are relatively large and amount to -9.43 kcal/mol (MP4(CP)(SDTQ)/TZP(2df,2pd)//MP2/TZP(2df,2pd)), -9.77 kcal/mol (DFT(CP)/TZP(2df,2pd))/DFT/TZP(2df,2pd)), -8.16 kcal/mol (CCSD(CP)/TZP(2df,2pd)//MP2/TZP(2df,2pd)), and -8.90 kcal/mol at the CCSD[T](CP)/TZP(2df,2pd) levels. The optimized planar molecular structure of the studied complex (MP2/TZP(2df,2pd) level) corresponds to the transition-state form (one imaginary, 66i cm<sup>-1</sup> harmonic vibrational frequency). A repulsive (2.07 kcal/mol) but minimum-energy planar structure of GeH<sub>2</sub>···OH<sub>2</sub> is predicted at the DFT/TZP(2df,2pd) level of approximation. However, single-point calculations at the planar DFT-optimized geometry reveal this complex to be slightly nonbounded (0.01 kcal/mol, CCSD-(CP)) or very weakly bounded (-0.17 kcal/mol, CCSD[T](CP)).

#### 1. Introduction

The matrix-isolation technique has been successfully applied to obtain a variety of new species and intermediates that contribute significantly to understanding principles of chemistry as well as atmospheric sciences as well. For example, an analysis of the recorded IR spectra of products from the photolytic reactions of silane/ozone and germane/ozone mixtures in low-temperature argon matrices revealed silicon- and germanium-containing analogs of well-known organic molecules such as methanol, formaldehyde, and formic and carbonic acids.<sup>1,2</sup> Also, molecular complexes, (e.g., HCl···HClCO, HBr····HBrCO, CO···(HCl)<sub>2</sub>, and CO···(HBr)<sub>2</sub><sup>3</sup> and phosgene complexes with F<sub>2</sub>, Cl<sub>2</sub>, and Br<sub>2</sub><sup>4,5</sup>) the intermediates and products of UV irradiation of halomethane/ozone mixtures have been detected. However, the IR investigations do not furnish detailed information on structure and energetics of these systems. For example, Whitnall and Andrews identified hydroxygermylene (HGeOH), germanone (H<sub>2</sub>GeO), and the GeH2...OH2 complex among products of matrix GeH4/O3 photolysis.<sup>2</sup> So performed a theoretical studies on molecular structure and relative stabilities of the isomeric HGeOH/H2-GeO<sup>6a</sup> and HGeSH/H<sub>2</sub>GeS<sup>6b</sup> species. More recently, Kapp et al. investigated structures and bonding in HXOH/H2XO (X = C, Si, Ge, Sn, and Pb) series by *ab initio* correlated methods.7

In the present paper, we present a quantum-mechanical study of the molecular structure and stability of the  $GeH_2\cdots OH_2$  complex.

#### 2. Computational Details

The equilibrium structures of the title system were optimized with the Gaussian92/94 set of programs<sup>8</sup> at the Hartree–Fock (HF), the second-order Møller-Plesset perturbation (MP2), and the density functional theory (DFT) with the exchange potential

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**Figure 1.** Molecular arrangement of the planar (top) and nonplanar (bottom)  $GeH_2\cdots OH_2$  complex.

of Becke and correlated functional of Lee, Yang, and Parr (Becke3LYP option in Gaussian terminology) correlated levels of theory within  $C_s$  (planar arrangement, see below) and  $C_1$  point-group symmetry of the complex. Additionally, single-point MP4(SDTQ) and coupled-cluster including single and double excitations (CCSD) and perturbative triple substitutions (CCSD[T]) calculations were performed at the MP2- and DFT-optimized geometries of the complex and its components. Two triple- $\zeta$ - quality basis sets augmented by sets of p and/or d (hydrogen atoms) and five-component d and/or seven-component f functions (Ge and O) were employed. The smaller basis set referred to as TZPH(2d,2p), consisted of the standard Pople's split-valence 6-311G with two sets of p-polarization functions for H, two sets of d-polarization functions for O,<sup>9</sup> and the

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TABLE 1: Selected Optimized Geometry Parameters for the GeH2...OH2 Complex and Its Components (in Parentheses)<sup>4</sup>

			TZPR (2df,2pd)			
	TZPH (2d,2p) nonplanar		nonplanar		planar	
	HF	MP2	MP2	DFT	MP2	DFT
		Bond Lengths				
O1-Ge2	2.286	2.214	2.214	2.268	3.687	3.720
Ge2-H3	1.582	1.579	1.577	1.607	1.561	1.591
		(1.570)	(1.569)	(1.597)		
Ge2-H4	1.574	1.569	1.567	1.595	1.563	1.590
O1-H5	0.942	0.961	0.962	0.963	0.957	0.964
		(0.957)	(0.957)	(0.961)		
O1-H6	0.942	0.961	0.961	0.963	0.960	0.961
		Bond Angles				
H3-Ge2-H4	94.0	92.6	92.4	91.4	92.9	91.9
		(91.7)	(91.7)	(90.9)		
H5-O1-H6	107.8	105.7	105.5	106.3	102.9	103.8
		(103.3)	(103.1)	(103.9)		
Ge2-H6-O1					176.6	172.3
		Dihedral Angles				
H6-O1-Ge2-H4	63.9	63.3	62.0	56.4		
H5-O1-Ge2-H4	173.6	177.0	175.4	170.5		
H5-O1-Ge2-H3	92.1	84.1	82.6	78.7		
H6-O1-Ge2-H3	-30.3	-29.6	-30.8	-35.4		
angle between GeH <sub>2</sub> and OH <sub>2</sub> plane	40.1	29.3	28.7	27.9		
angle between O1–H6 and $GeH_2$ plane	27.5	23.0	22.6	22.8		
angle between O1-H5 and GeH <sub>2</sub> plane	17.3	11.6	11.3	10.0		
O distance from GeH <sub>2</sub> plane	2.276	2.203	2.202	2.255		

<sup>a</sup> Distances in Å, angles in degrees.

partially decontracted [433111/43111/4\*] Huzinaga's basis supplemented with two d-polarization functions<sup>10</sup> for Ge. A larger basis set was obtained by saturation of the previous basis set with additional d-polarization (H) and f-polarization sets (O). For Ge, the 6-311G-like basis set of Curtiss et al.<sup>11</sup> augmented with one set of d- and one set of f-polarization functions was used. This basis set will be refered to as the TZPR(2df,2pd). Spherical five-component d and seven-component f polarization functions were used. Calculated interaction energies were corrected for the basis set superposition error using the full Boys-Bernardi counterpoise (CP) scheme.<sup>12</sup> The interaction energies were partitioned into the SCF ( $\Delta E^{SCF}$ ) and second- $(\Delta E^2)$ , third- $(\Delta E^3)$ , and fourth-order  $(\Delta E(\text{SDTQ})^4)$  intermolecular Møller–Plesset perturbation theory components.  $\Delta E^{SCF}$ was divided into Heitler-London first-order electrostatic and exchange ( $\Delta E^{(1)}$ ), and mutual induction (deformation),  $\Delta E^{\text{SCF}}$ -(def) terms:  $\Delta E^{\text{SCF}} = \Delta E^{(1)} + \Delta E^{\text{SCF}}$ (def) (for example, see: ref 13). Pure cartesian d- (six-component) and f- (10component) polarization functions were used.<sup>14</sup> Frozen core (1s2s2p3s3p for Ge and 1s for O atomic orbitals) approximation was adopted at the MP2, CCSD, CCSD[T], and MP4 levels of theory.

### 3. Results and Discussion

Optimized molecular arrangements assuming planar ( $C_s$ ) and unconstrained  $C_1$  symmetry of the GeH<sub>2</sub>····OH<sub>2</sub> system are displayed schematically in Figure 1. As a possible structure for the title complex, Whitnall and Andrews<sup>2</sup> suggested coplanar orientation of the H<sub>2</sub>O and GeH<sub>2</sub> subunits with the Ge···O contact. However, depending on the applied level of approximation, the stationary-point structures for the planar H<sub>2</sub>O····GeH<sub>2</sub> complex correspond to a second-order transition form at the MP2/TZPH(2d,2p) (two imaginary harmonic vibrational frequencies of 147i and 71i cm<sup>-1</sup>), a first-order transition form at the MP2/TZPR(2df,2pd) (one imaginary frequency of 66i cm<sup>-1</sup>), and a minimum-energy structure at the DFT/TZPR(2df,2pd) levels of theory. Profoundly nonplanar global-minimum geometries of the complex are predicted at all

applied levels of theory. The H<sub>2</sub>O and GeH<sub>2</sub> subunits form two nonparallel and twisted planes (Figure 1 and Table 1). The HF/TZPH(2d,2p)- and DFT/TZPR(2df,2pd)-optimized structures of the title complex correspond to more separated GeH<sub>2</sub> and H<sub>2</sub>O subunits. The O–Ge intermolecular distance (2.286 Å at the HF/TZPH(2d,2p) and 2.268 Å at the DFT/TZPR(2df,2pd)) and oxygen atom distance from the GeH<sub>2</sub> plane (2.276 Å at HF/TZPH(2d,2p) and 2.255 Å at DFT/TZPR(2df,2pd)) are noticeably larger than the corresponding values predicted at the MP2/TZPH(2d,2p) and the MP2/TZPR(2df,2pd) (Ge-O, 2.214 Å and 2.214 Å; O-GeH<sub>2</sub> plane, 2.203 and 2.202 Å, respectively) levels of theory. Intermolecular angles for  $GeH_2 \cdots OH_2$ predicted at the Hartree-Fock level correspond to the most nonparallel and twisted structure. The interplanar angle amounts to 40.1°, and the angles formed by the O-H bonds with the GeH<sub>2</sub> plane (27.5° and 17.3°) are larger than values optimized at correlated levels (interplanar angle, 29.3°, 28.7, and 27.9; O-H6 angle with GeH<sub>2</sub>, 23.0°, 22.6°, and 22.8°; O-H5 angle with GeH<sub>2</sub>, 11.6°, 11.3°, and 10.0° at the MP2/TZPH(2p,2d), MP2/TZPR(2df,2pd), and DFT/TZPR(2df,2pd), respectively). At all levels, the typical geometric changes of GeH<sub>2</sub> and H<sub>2</sub>O upon nonplanar complex formation are observed (i.e., an elongation of bond lengths, (ca. 0.01 Å) and an increase of bond angles  $(0.5^{\circ}-3^{\circ})$ ). Smaller distortions of the subunits that form the planar complex structure are predicted (Table 1).

The investigated complex is relatively strong. Listed in Table 2 are the calculated total electronic, zero-point vibrational, and interaction energies ( $\Delta E$ ) for the nonplanar and planar GeH<sub>2</sub>···OH<sub>2</sub> complexes and its component molecules supplemented by values of dipole moments and dipole polarizabilities. The SCF interaction energy component does not differ significantly at the MP2-optimized ( $\Delta E^{\text{SCF}}(\text{CP}) = -5.40$  kcal/mol (TZPH(2d,2p)) and -5.70 kcal/mol (TZPH(2d,2pd))) or HF-optimized ( $\Delta E^{\text{SCF}}(\text{CP}) = -6.23$  kcal/mol (TZPH(2d,2p)) geometries of the nonplanar GeH<sub>2</sub>···OH<sub>2</sub>. These considerably large attractive values of  $\Delta E^{\text{SCF}}$  originate mainly from the contribution of the deformation term. Whereas GeH<sub>2</sub> has practically no permanent dipole moment (0.05 D at HF/TZPH-

TABLE 2: Calculated Total Electronic (*E*), Zero-Point Vibrational (ZPE, in au), Interaction Energies ( $\Delta E$ , in kcal/mol), Dipole Moments ( $\mu$ , in D), and Dipole Polarizabilities ( $\alpha$ , in au) for the GeH<sub>2</sub>…OH<sub>2</sub> Complex and Its Component Molecules

	TZPH (2d,2p)		TZPR (2df,2pd)					
				nonplanar			planar	
	MPn	HF	MPn	$\mathrm{CCSD}^a$	DFT	MPn	DFT	$\mathrm{CCSD}^b$
$\overline{E(\text{GeH}_2 \cdots \text{OH}_2)}$	-2151.031 170 <sup>c</sup>	-2150.647 778	-2153.035 514 <sup>c</sup>	-2153.050 379	-2154.640 625	-2153.010 228 <sup>c</sup>	-2154.617236	-2153.028124
	$-2151.049\ 035^{d}$		$-2153.038\ 372^{d}$	-2153.067 471		-2153.015 692d		-2153.043 996
	$-2051.070\ 820^{e}$		$-2153.071\ 146^{e}$			$-2153.046957^{e}$		
ZPE	0.039 074	0.040396	0.038831		0.037 523	0.034 808	0.033927	
μ	3.70	3.76	3.78		3.62	2.24	2.03	
$E(H_2O)$	$-76.286\ 480^{\circ}$	-76.051 109	-76.310 157°	-76.316 757	$-76.453\ 418$			-76.316 696
	76.290 246 <sup>d</sup>		$-76.315\ 045^{d}$	-76.316 757				-76.323 997
	$-76.299\ 464^{e}$		$-74.324\ 685^{e}$					
ZPE	0.021 755	0.023 153	0.021 800		0.021 339			
μ	1.90	1.96			1.87			
$\alpha, \alpha_{\perp}$	6.74, 8.14	6.27, 7.64	6.79, 8.16		6.93, 7.08			
$E(GeH_2)$	$-2074.714534^{\circ}$	-2074.576842	$-2076.696\ 001^{c}$	-2076.708 343	-2078.160 600			-2076.708 238
,	$-2074.731918^{d}$		$-2076.697\ 065^{d}$	-2076.716 493				-2076.716 438
	$-2074.742~727^{e}$		-2076.718 440e					
ZPE	0.011 494	0.011 635	0.011 353		0.010 732			
и		0.05	0.00		0.28			
ά, αι	34.13, 37.40	34.76, 36.93	30.61, 35.63		31.62, 36.22			
$E(H_2O)$ dim	$-76.293096^{\circ}$	-76.054 583	$-76.316\ 100^{\circ}$	-76.322010	$-76.458\ 810$			
/	$-76.295\ 800^{d}$		$-76.320055^{d}$	-76.329 693		$-76.315926^{d}$		-76.324 992
	$-76.306\ 112^{e}$		$-76.330~597^{e}$			$-76.325\ 750^{e}$		
E(GeH <sub>2</sub> ) dim	$-2074.717\ 520^{\circ}$	-2074.577 654	$-2076.697\ 411^{c}$	-2076.709 691	-2078.160794	$-2076.696\ 326^{\circ}$	-2078.160 596	-2076.708 664
· _/	$-2074.731918^{d}$		$-2076.698 418^{d}$	-2076.717912		$-2076.697\ 339^{d}$		-2076.716882
	$-2074.745~786^{e}$		$-2076.719\ 848^{e}$			$-2076.718724^{e}$		
$\Delta E^{\rm SCF}$	-8.34	-8.92	-7.93		-13.27		1.40	
$\Delta E^{\text{SCF}}(\text{CP})$	-5.40	-6.23	-5.70		-9.77		2.07	
$\Delta E^2$	-6.93		-6.93			1.01		
$\Delta E^2(CP)$	-3.84		-4.54			1.88		
$\Delta E^{3}$	2.06		-0.94					
$\Delta E^3$ (CP)	-2.33		-1.32					
$\Delta E^4$	-1.10		-0.10					
$\Delta E^4(CP)$	3.35		-0.54					
$\Delta E(4)$	-14.31		-14.02			1.16		
$\Delta E(4)(CP)$	-8.22		-9.43			2.05		
CCSD				-12.30				-0.84
CCSD[T]				-13.34				-1.07
CCSD(CP)				-8.16				0.01
CCSD[T](CP)				-8.90				-0.17

<sup>*a*</sup> At the MP2-optimized geometry. The first entry corresponds to CCSD, the second to CCSD[T]. Interaction energy values are corrected for ZPE(MP2). <sup>*b*</sup> At the DFT-optimized geometry. The first entry corresponds to CCSD, the second to CCSD[T]. Interaction energy values are corrected for ZPE(DFT). <sup>*c*</sup> MP2 values at the MP2-optimized geometry. <sup>*d*</sup> MP3 values at the MP2-optimized geometry. <sup>*e*</sup> MP4(SDTQ) values at the MP2-optimized geometry.

(2d,2p) and 0.03 D at HF/TZPR(2df,2pd)//MP2 /TZPR(2df,-2pd) level), high values of the dipole polarizability are predicted: the average (34.76 au at the HF/TZPH(2d,2p) and 30.61 au at HF/TZPR(2df,2pd)//MP2/TZPR(2df,2pd) level) and its perpendicular component, (36.93 au and 35.63 au, respectively) are especially important due to the nonplanar orientation of the GeH<sub>2</sub> and H<sub>2</sub>O planes. At the MP2/6-311G(2df,2pd)optimized geometry, the following values of the SCF interaction energy components were obtained:  $\Delta E^{(1)}(CP) = 9.49, \Delta E^{SCF}$ -(def, CP) = -18.73, and  $\Delta E^{SCF}(CP) = -9.24$  kcal/mol, respectively.<sup>15</sup> The second-order component ( $\Delta E^{(2)}(CP)$ ), containing dispersion and higher induction terms, contributes very significantly to the total interaction energy and is of comparable magnitude to the SCF component. It constitutes 71% (TZPH-(2d,2p)) and even 80% of  $\Delta E^{\text{SCF}}(\text{CP})$  with the TZPR(2df,2pd) basis set. The higher-order terms reveal opposite basis-set dependencies. Calculated with the TZPH(2d,2p) basis set the  $\Delta E^{(3)}(CP)$  and  $\Delta E^{(4)}(CP)$  contributions are large and the  $\Delta E^{(4)}$ -(CP) correction is noticeably larger (3.35 kcal/mol) than  $\Delta E^{(3)}$ -(CP) (-2.33 kcal/mol). Moreover, the basis-set extension effect (i.e., the difference between values obtained with monomerand dimer-centered ( $\Delta E^{(x)}(CP)$ ) basis sets) is enormously large (-2.33 vs 2.06 kcal/mol and 3.35 vs -1.10 kcal/mol in thirdand fourth-order, respectively), and this suggests that the TZPH-(2d,2p) basis set is not saturated sufficiently to describe intramolecular correlation effects. A better convergence of the MP perturbation theory is achieved with the TZPR(2df,2pd) basis set: absolute values of the second-, third-, and fourth-order corrections decrease monotonically, and the total interaction energy  $\Delta E4(CP) = -9.43$  kcal/mol differs insignificantly from the value of -9.78 kcal/mol calculated at the CCSD[T]-(CP) level value.

The good performance of the DFT method should be noted, and the predicted total interaction energies of DFT(CP)/TZPR-(2df, 2pd) = -9.77, MP4(SDTQ)(CP)/TZPR(2df, 2pd) = -9.43, and CCSD[T](CP) = -8.90 kcal/mol do not differ significantly (Table 2). As was mentioned above, the planar conformation of GeH2...OH2 is predicted to be a minimum-energy isomer only at the DFT/TZPR(2df,2pd) level of theory; however, the calculated interaction energy corresponds to a repulsive ( $\Delta E(CP)$ ) = 2.07 kcal/mol) structure. On the other hand, the  $C_s$ conformation optimized at the MP2/TZPR(2df,2pd) level is a transition-state form (one imaginary harmonic frequency of 66i  $cm^{-1}$ , Table 3) that is repulsive at the both MP2 (1.88 kcal/ mol) and MP4(SDTQ)(CP)/TZPR(2df,2pd)//MP2/TZPR(2df,-2pd) (2.05 kcal/mol). Using the DFT-optmized planar geometry of GeH<sub>2</sub>…OH<sub>2</sub>, nonbonded (0.01 kcal/mol) or very weakly attractive (-0.17 kcal/mol) structures are predicted at the CCSD-(CP)/TZPR(2df,2pd) and CCSD[T](CP)/TZPR(2df,2pd) levels of approximation.

# TABLE 3: Calculated Harmonic Vibrational Frequencies (Unscaled, in cm<sup>-1</sup>) and Absolute IR Intensities (in km/mol) for GeH<sub>2</sub>···<sup>16</sup>OH<sub>2</sub> at the MP2/TZPR(2df,2pd) and DFT/TZPR(2df,2pd) Levels

MP2/DFT/exptl. <sup>a</sup> nonplanar minimum energy form <sup>a</sup>	MP2/DFT planar transition form
A: 112 (37), 116 (38) <sup>b</sup> rotations of GeH <sub>2</sub> and OH <sub>2</sub> around Ge $-$ O axis A: 274 (42), 249 (40) Ge $-$ O stretching	A": 66i (89), 18 (94) butterfly-like asym motion of GeH <sub>2</sub> and OH <sub>2</sub> A": 45 (5), 77 (24) libration of OH <sub>2</sub> against GeH <sub>2</sub>
A: $430 (178), 400 (171)$ butterfly-like asym motion of OH <sub>2</sub> and GeH	A': $67(9)$ , $82(0)$ notation of GeH <sub>2</sub> and OH <sub>2</sub> against each other, G=O stratching
A: 491 (11), 462 (11) butterfly-like asym motion of $OH_2$ and GeH	A': 87 (0), 155 (30) libration of $OH_2$ against GeH <sub>2</sub> , rotations of
A: 665 (48), 624 (50) butterfly-like sym motion of $OH_2$ and GeH	$A'$ : 218 (59), 213 (31) libration of $OH_2$ against $GeH_2$ rotations of
A: 751 (33), 702 (27) butterfly-like sym motion of OH <sub>2</sub> and GeH <sub>2</sub>	GeH <sub>2</sub> and OH <sub>2</sub> around Ge-O axis A": 329 (47), 298 (29) butterfly-like sym motion of GeH <sub>2</sub> and OH <sub>2</sub>
A: 956 (94), 935 (81), 897.8 GeH <sub>2</sub> bending	A': 967 (90), 945 (82) GeH <sub>2</sub> bending
A: 1644 (71), 1629 (67), 1586.1 OH <sub>2</sub> bending	A': 1675 (52), 1653 (72) OH <sub>2</sub> bending
A: 1949 (427), 1812 (446), 1777.2 GeH stretching	A': 2030 (299), 1902 (320) GeH <sub>2</sub> stretching sym
A: 1997 (307), 1809 (392), 1813.0 Gen stretching A: 3830 (80) 3788 (65) 3597 / OH, stretching sym	A : 2039 (301), 1914 (303) GeH <sub>2</sub> stretching asym $\Delta'$ : 3847 (131) 3758 (137) OH <sub>2</sub> stretching sym
A: $3945 (134), 3885 (118), 3686.0 \text{ OH}_2 \text{ stretching asym}$	A': 3974 (117), 3876 (88) OH <sub>2</sub> stretching asym

<sup>*a*</sup> In Ar matrix, ref 2. <sup>*b*</sup> For all entries, symmetry: frequency (integration).

TABLE 4: Calculated (MP2/TZPR(2df,2pd) Vibrational Frequencies upon Isotopic Substitution for GeH<sub>2</sub> and H<sub>2</sub>O Subunits of the GeH<sub>2</sub>…OH<sub>2</sub> Complex<sup>*a*</sup>

mode	GeH <sub>2</sub>	GeHD	$GeD_2$	
$\delta$ (H-Ge-H) $\nu$ (Ge-H)	956 (897.8) 1949, 1997 (1777.2, 1813.6)	833, 831 1388, 1422, 1948, 1997 (1287.0, 1307.4, 1782.0, 1811.6)	690 (646.4) 1387, 1442 (1281.6, 1308.9)	
	H <sub>2</sub> O s			
mode	H <sub>2</sub> <sup>16</sup> O/H <sub>2</sub> <sup>18</sup> O	HDO	$D_2O$	
$\delta$ (H–O–H) $\nu$ <sub>s</sub> (O–H) $\nu$ <sub>as</sub> (O–H)	1644/1638 (1586.1/1580.1) 3830/3822 (3597.4/3590.3) 3945/3929 (3686.0/3672.5)	1443, 1445 (1398.4) 2822, 2825 (2682.3) 3888, 3892 (3637.3)	1204 (1173.6) 2760 (2627.8) 2892 (2738.5)	

<sup>*a*</sup> In cm<sup>-1</sup>. Experimental data in parentheses (Ar matrix, ref 2).

Calculated harmonic vibrational frequencies and IR intensities for the GeH<sub>2</sub>…OH<sub>2</sub> complex are listed in Table 3. Lowfrequency vibrations corresponding to intermolecular motions for planar and nonplanar forms differ substantially due to their different molecular symmetries. Intramolecular stretching and bending modes of component subunits predicted at both the MP2 and the DFT/TZPR(2df,2pd) levels for nonplanar conformation are lower and closer to experimental data than corresponding values calculated for the planar GeH2...OH2 (Table 3). As usual, the calculated (unscaled) harmonic vibrational frequencies, especially for higher-lying stretching vibrations, are larger than the corresponding measured anharmonic values. However, measured and calculated shifts of frequencies for isotopic nonplanar GeH2····OH2 systems are very close. For example, observed diagnostic values of the <sup>18</sup>O red shift of the O-H stretching modes (7.1 and 13.4 cm<sup>-1</sup>) and the MP2/TZPR-(2df,2pd)-predicted (8 and 16 cm<sup>-1</sup>) values agree very well (Table 4).

## 4. Conclusions

Molecular parameters and interaction energy of the GeH<sub>2</sub>···OH<sub>2</sub> complex detected by Whitnall and Andrews in argon matrices<sup>2</sup> were calculated using the Hartree–Fock and post-Hartree–Fock quantum-mechanical methods. The obtained results strongly support a nonplanar geometry of the complex; its nonplanar structure corresponds to a relatively strongly bonded system for which the interaction energy insignificantly varies at the applied correlated levels of theory (*e.g.*,  $\Delta E(CP) = -9.43$  kcal/mol at the MP4(SDTQ)(CP)/TZPR-(2df,2pd)//MP2/TZPR(2df,2pd), -9.77 kcal/mol at the DFT-(CP)/TZPR(2df,2pd)//DFT/TZPR(2df,2pd), and -8.90 kcal/mol at the CCSD[T](CP)/TZPR(2df,2pd)//MP2(2df,2pd)/TZPR-

(2df,2pd) levels). The optimized molecular structure obtained assuming planarity corresponds to a transition-state repulsive form (1.88 kcal/mol) at the MP2(CP)/TZPR(2df,2pd)//MP2/TZPR(2df,2pd) and 2.05 kcal/mol at the MP4(SDTQ)(CP)/TZPR(2df,2pd)//MP2/TZPR(2df,2pd) levels). The planar structure of the title complex is predicted to be a minimum-energy but also repulsive (2.07 kcal/mol) conformation at the DFT-(CP)/TZPR(2df,2pd) level. Results of coupled-cluster calculations based on the DFT-optimized geometry are not conclusive: 0.01 kcal/mol at the CCSD(CP) and -0.17 kcal/mol at the CCSD[T](CP).

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#### **References and Notes**

- (1) Whitnall, R.; Andrews, L. J. Phys. Chem. 1985, 89, 3261.
- (2) Whitnall, R.; Andrews, L. J. Phys. Chem. 1990, 94, 2351.
- (3) Lugez, C.; Schriver, A.; Schriver-Mazzuoli, L.; Lasson, E.; Nielsen, C. J. J. Phys. Chem. **1993**, *97*, 11617.
- (4) Abdelaoui, O.; Schriver, L.; Schriver, A. J. Mol. Struct. 1992, 268, 335.

(5) Bouteiller, Y.; Abdelaoui, O.; Schriver, A.; Schriver-Mazzuoli, L. J. Chem. Phys. **1995**, 102, 1731.

(6) (a) So, S. P. J. Phys.Chem. **1991**, 95, 10658. (b) So, S. P. J. Phys. Chem. **1993**, 97, 4643.

(7) Kapp, J.; Remko, M.; Schleyer P. v. R. J. Chem. Phys. 1996, 118, 5745.

(8) (a) Frisch, M. J.; Trucks, G. W.; Head-Gordon, M.; Gill, P. M. W.; Wong, M. W.; Foresman, J. B.; Johnson, B. G.; Schelgel, H. B.; Robb, M.; Replogle, E. S.; Gomperts, R.; Anders, J. L.; Raghavachari, K.; Binkley, J. S.; Gonzales, C.; Martin, R. I.; Fox, D. J.; Defrees, D. J.; Baker, J.; Stewart, J. J. P.; Pople, J. A.; *Gaussian 92*, Revision A; Gaussian, Inc.: Pittsburgh, PA, 1992. (b) Frisch, M. J.; Truck, G. W.; Schlegel, H. B.; Gill, P. M. W.; Johnson B. G.; Robb, M. A.; Cheeseman, J. R.; Keith, T.; Petersson, G. A.; Montgomery, J. A.; Raghavachari, K.; Al-Laham, M. A.; Zakrzewski, V. G.; Ortiz, J. V.; Foresman, J. B.; Peng, C. Y.; Ayala, P. Y.; Chen, W.; Wong, M. W.; Andres, J. L.; Replogle, E. S.; Gomperts, R.; Martin, R. L.; Fox, D. J.; Binkley, J. S.; Defrees, D. J.; Baker, J.; Stewart, J. J. P.; Head-Gordon, M.; Gonzalez, C.; Pople, J. A. *GAUSSIAN 94*, Revision B.3; Gaussian, Inc.: Pittsburgh, PA, 1995.

(9) (a) Krishnan, R.; Binkley, J. S.; Seeger, R.; Pople, J. A. J. Chem. Phys. **1980**, 72, 650. (b) Frisch, M. J.; Pople, J. A.; Binkley, J. S. J. Chem. Phys. **1984**, 80, 3265.

(10) Huzinaga, S.; Andzelm, J.; Klobukowski, M.; Radzio-Andzelm, E.; Sakai, Y.; Tatewaki, H. *Gaussian Basis Sets for Molecular Orbital Calculations;* Elsevier: New York, 1984.

(11) Curtiss, L. A.; McGrath, M. P.; Blaudeau, J.-P.; Davis, N.; Binning, R. C., Jr.; Radom, L. J. Chem. Phys. **1995**, 103, 6104.

(12) Boys, S. F.; Bernardi, F. Mol. Phys. 1970, 19, 553.

(13) Chatasiński, G.; Szczęśniak, M. M. Mol. Phys. 1988, 63, 205.

(14) Slightly modified version of GAMESS. See: Schmidt, M. W.;

Baldridge, K. K.; Boatz, J. A.; Jensen, J. H.; Koseki, S.; Gordon, M. S.; Nguyen, K. A.; Windus, T. L.; Elbert, S. T. *QCPE* **1990**, *10*, 52.

(15) A discrepancy between cited above  $\Delta E^{\text{SCF}}(\text{CP}) = -5.70 \text{ kcal/mol}$  (HF/TZPR(2df,2pd))/MP2/TZPR(2df,2pd)) and here (-9.24 kcal/mol), which are the results from different size of the basis sets used: spherical (5d and 7f, Gaussian 92/94) and pure cartesian (6d and 10f, GAMESS) Gaussian-type orbitals.